

Fits to ν -CCQE Using Model-Independent Axial Form Factor Parameterization

Ab Initio Approaches

Aaron Meyer (asmeyer2012@uchicago.edu)

University of Chicago/Fermilab

July 15, 2015

Motivation

Neutrino physics needs a better understanding of axial form factor:

- **Model-dependent shape parameterization** introduces systematic uncertainties and underestimates errors
- **Nuclear effects entangled** with nucleon cross sections
- Measurement of oscillation parameters depends on **nuclear models** and **nucleon-level form factors**

Why Do We Still Need Better Theory?

Neutrino physics uses near detector/far detector paradigm,
measures number distribution:

$$\frac{N_{\text{CCQE, near}}(E_\nu)}{N_{\text{CCQE, far}}(E_\nu)} = \frac{\phi_{\text{near}}(E_\nu) \sigma_{\text{CCQE}}(E_\nu) \epsilon_{\text{near}}}{\phi_{\text{far}}(E_\nu) \sigma_{\text{CCQE}}(E_\nu) \epsilon_{\text{far}}}$$

Problems:

- ϵ depends on near/far detector technology
- σ depends on nuclear models/nuclear target at near/far
- ϕ depends on beam angular distribution
→ near/far detector sample different energy distributions

Why Do We Still Need Better Theory?

Neutrino physics uses near detector/far detector paradigm, measures number distribution:

$$\frac{N_{\text{CCQE, near}}(E_\nu)}{N_{\text{CCQE, far}}(E_\nu)} = \frac{\phi_{\text{near}}(E_\nu) \sigma_{\text{CCQE}}(E_\nu) \epsilon_{\text{near}}}{\phi_{\text{far}}(E_\nu) \sigma_{\text{CCQE}}(E_\nu) \epsilon_{\text{far}}}$$

More Problems:

- σ is modified by nuclear and radiative corrections
- Effects of corrections removed by studying modification of N with Monte Carlo
- Monte Carlo uses σ as input
- σ calculated by measuring N

Degenerate uncertainties $N \rightarrow \text{MC} \rightarrow \sigma \rightarrow N$

Why Do We Still Need Better Theory?

Neutrino physics uses near detector/far detector paradigm,
measures number distribution:

$$\frac{N_{\text{CCQE, near}}(E_\nu)}{N_{\text{CCQE, far}}(E_\nu)} = \frac{\phi_{\text{near}}(E_\nu) \sigma_{\text{CCQE}}(E_\nu) \epsilon_{\text{near}}}{\phi_{\text{far}}(E_\nu) \sigma_{\text{CCQE}}(E_\nu) \epsilon_{\text{far}}}$$

Even More Problems:

- Model for σ constructed from single-nucleon cross section
- single-nucleon cross section constrained by assuming a model for σ

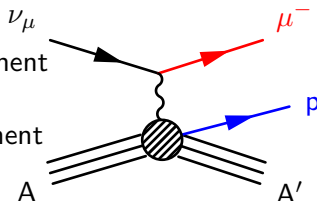
Degenerate uncertainties $\sigma_A \rightarrow \sigma_N \rightarrow \sigma_A$

Nuclear Effects

Nuclear effects not well understood

→ Models which are best for one measurement
are worst for another

Need to break F_A /nuclear model entanglement



(assumed $m_A = 0.99$ GeV, reference hyperlinks online)

NuWro Model (χ^2 /DOF)	RFG [GENIE]	RFG+ TEM	assorted others
leptonic(rate)	3.5	2.4	2.8-3.7
leptonic(shape)	4.1	1.7	2.1-3.8
hadronic(rate)	1.7[1.2]	3.9	1.9-3.7
hadronic(shape)	3.3[1.8]	5.8	3.6-4.8

Discrepancies in the Axial-Vector Form Factor

Most analyses assume the “Dipole form factor”:

$$F_A^{\text{dipole}}(q^2) = g_A \frac{1}{\left(1 - \frac{q^2}{m_A^2}\right)^2}$$

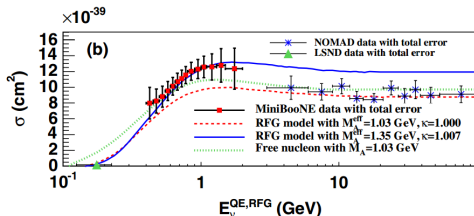
Dipole is an ansatz:

unmotivated in interesting energy range

→ **uncontrolled systematics** and **underestimated uncertainties**

Essential to replace ansatz with

model-independent parameterization



MiniBooNE Collab., PHYS REV D 81, 092005 (2010)

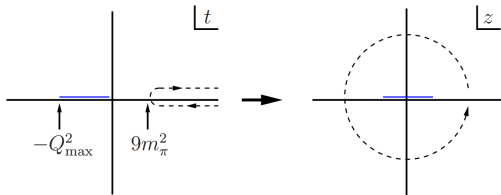
z-Expansion

The z-Expansion ([Bhattacharya, Hill, Paz arXiv:1108.0423 \[hep-ph\]](#)) is a conformal mapping which takes the kinematically allowed region ($t \leq 0$) to within $z = \pm 1$

→ For reference, later plots will have $|z_{\max}| = 0.28$

$$t = q^2 = -Q^2 \quad t_c = 9m_\pi^2$$
$$z(t; t_0, t_c) = \frac{\sqrt{t_c - t} - \sqrt{t_c - t_0}}{\sqrt{t_c - t} + \sqrt{t_c - t_0}}$$

$$F_A(z) = \sum_{n=0}^{\infty} a_n z^n$$



z-Expansion implemented in [GENIE](#), to be released soon [autumn]

Advantages of z-Expansion

z-Expansion is a **model-independent** description of the axial form factor

- Motivated by analyticity arguments
- Provides a prescription for introducing more parameters as data improves
- Allows quantification of systematic errors

From meson (baryon) semileptonic decays, only a few expansion coefficients necessary to accurately represent data

- Coefficient falloff required by perturbative QCD
- For general analysis, see [Hill \[arXiv:hep-ph/0606023\]](#)
- For recent $|V_{ub}|$ determination, see [Fermilab/MILC \[arXiv:1503.07839\]](#)
- For recent $|V_{ub}|/|V_{cb}|$ determination, see [LHCb \[arXiv:1504.01568\]](#)

Evaluation of Fits

Process:

- Fit to increasing k_{\max} until adding new parameters no longer contributes appreciably to error/shape (z^k small)

What to Expect:

- Errors monotonically increase with more parameters
- Higher order coefficients alter fit less than lower order
 - Can cut off at finite k_{\max} with marginal impact to fit
 - Data indicates how many parameters should be used
 - Truncation error is a systematic
- Coefficients $O(1)$, decreasing amplitude as k increases
- Expect shape to fit data, no other requirements on shape

Deuterium Fitting

with Richard Hill, Rik Gran, Minerba Betancourt

Fitting done on deuterium bubble chamber data
(controlled nuclear effects)

Three datasets (reference hyperlinks online):

- [ANL 1982](#): 1737 events, 0.5GeV [peak]
- [BNL 1981](#): 1138 events, 1.6 GeV [average]
- [FNAL 1983](#): 362 events, 20 GeV [peak], 27 GeV [average]

PRELIMINARY shape-only fits to QE differential cross section data

Results propagated to single nucleon QE total cross section

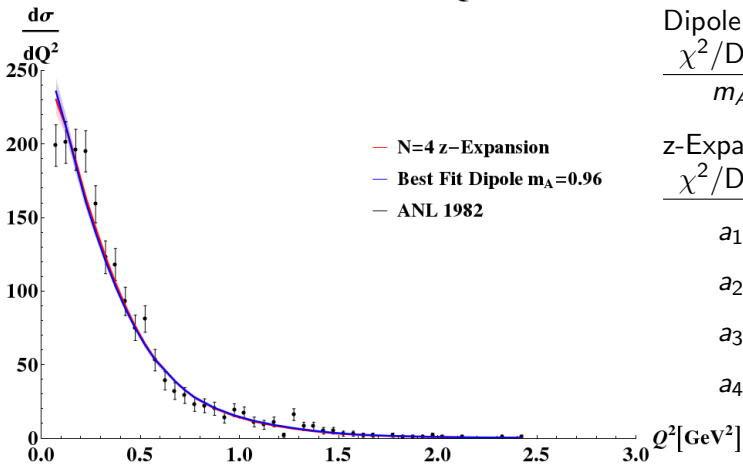
Gaussian priors used on z-Expansion coefficients:

if ($k \leq 5$) $\sigma_k = 5$, else $\sigma_k = 25/k$

Sum rule applied to ensure $F_A \sim 1/Q^4$ as $Q^2 \rightarrow \infty$

Deuterium Fitting Results

PRELIMINARY – ANL 1982 QE Diff xsec

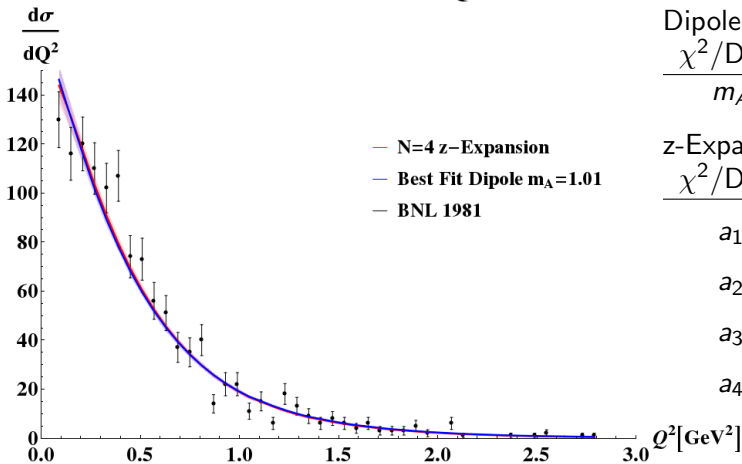


Dipole:	
χ^2/DOF	55.3/39
m_A	0.96(5)

z-Expansion:	
χ^2/DOF	54.5/36
a_1	$2.43^{+0.13}_{-0.12}$
a_2	$0.66^{+1.09}_{-1.05}$
a_3	$-7.18^{+2.80}_{-2.87}$
a_4	$4.26^{+3.46}_{-3.45}$

Deuterium Fitting Results

PRELIMINARY – BNL 1981 QE Diff xsec

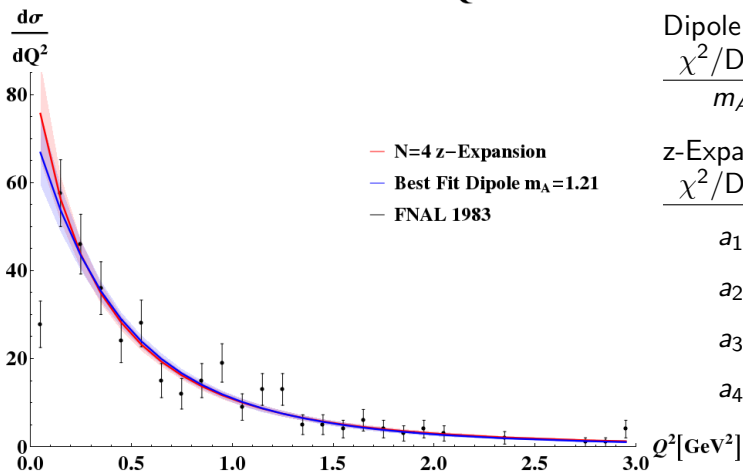


Dipole:	
χ^2/DOF	53.3/36
m_A	1.01(5)

z-Expansion:	
χ^2/DOF	53.9/33
a_1	$2.30^{+0.12}_{-0.12}$
a_2	$0.88^{+1.22}_{-1.26}$
a_3	$-6.55^{+3.01}_{-3.07}$
a_4	$3.05^{+3.49}_{-3.49}$

Deuterium Fitting Results

PRELIMINARY – FNAL 1983 QE Diff xsec



Dipole:

χ^2/DOF	15.9/21
m_A	1.21(12)

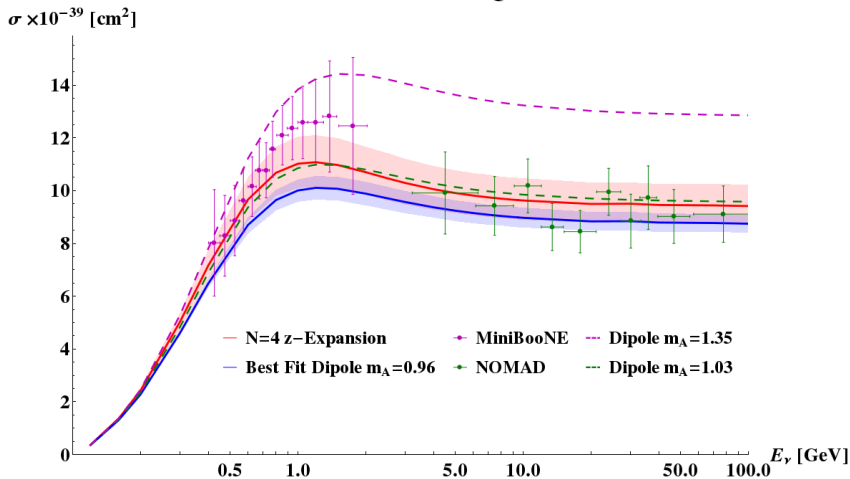
z-Expansion:

χ^2/DOF	15.3/18
a_1	$1.95^{+0.16}_{-0.16}$
a_2	$-1.32^{+1.81}_{-1.80}$
a_3	$0.03^{+3.67}_{-3.67}$
a_4	$-0.11^{+3.56}_{-3.56}$

Deuterium Fitting Results

- Quoted NOMAD/MiniBooNE σ are $\frac{1}{6}\sigma_{\text{carbon}}$
- Experiments use different definitions of CCQE
- Dipole guide lines (dashed) are nucleon-level cross section

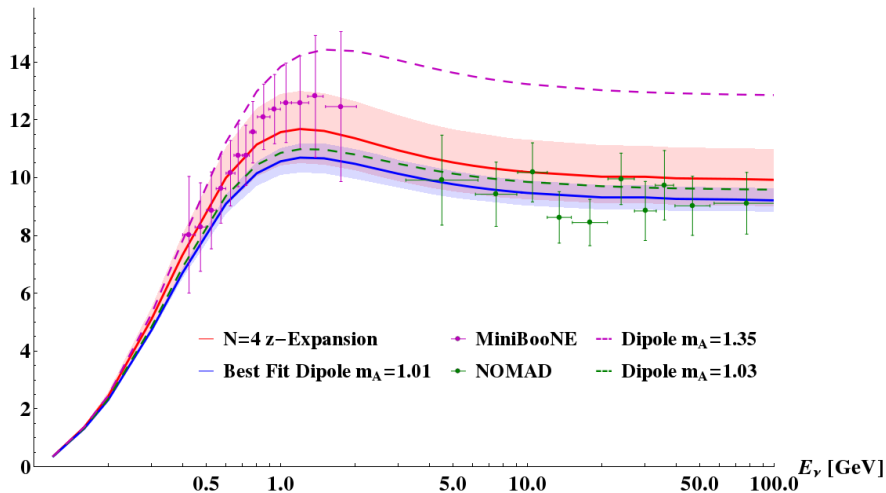
PRELIMINARY – ANL 1982 QE Cross Section



Deuterium Fitting Results

PRELIMINARY – BNL 1981 QE Cross Section

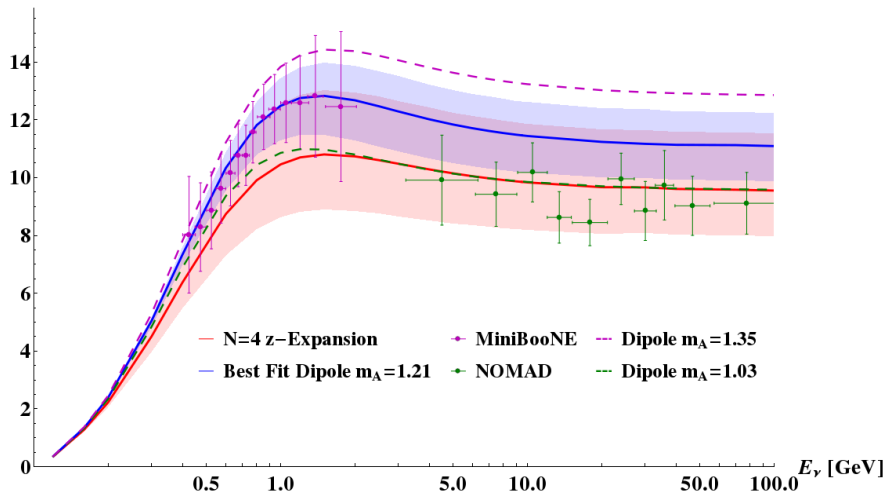
$\sigma \times 10^{-39} [\text{cm}^2]$



Deuterium Fitting Results

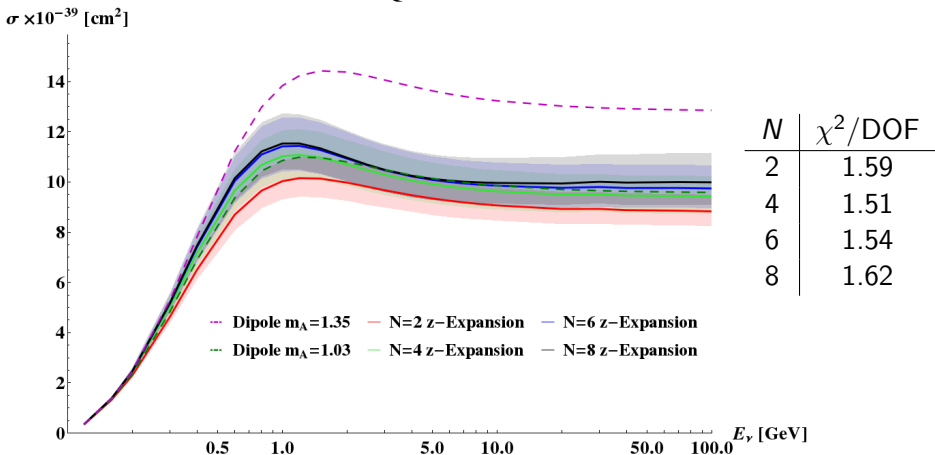
PRELIMINARY – FNAL 1983 QE Cross Section

$\sigma \times 10^{-39} [\text{cm}^2]$



Deuterium Fitting Results

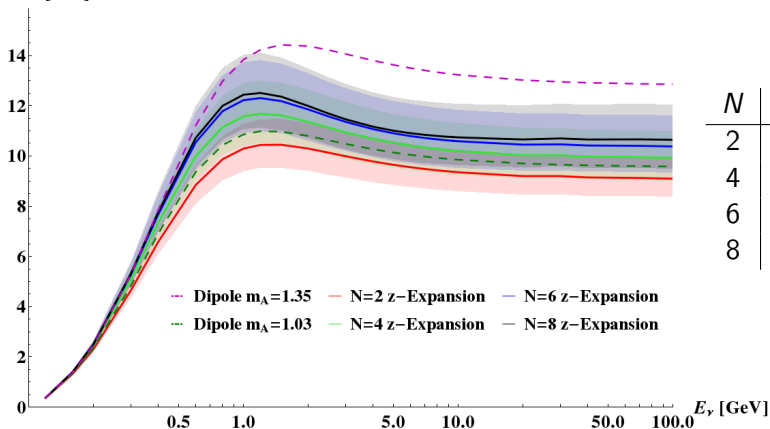
ANL 1982 Total QE Cross Section



Deuterium Fitting Results

BNL 1981 Total QE Cross Section

$\sigma \times 10^{-39} [\text{cm}^2]$

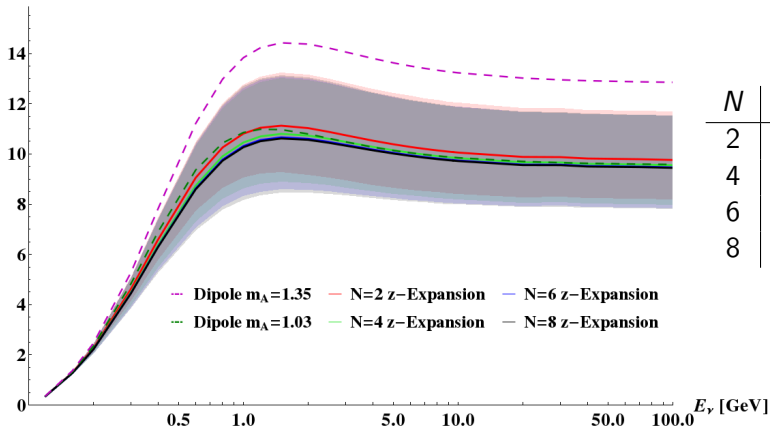


N	χ^2/DOF
2	1.67
4	1.63
6	1.66
8	1.74

Deuterium Fitting Results

FNAL 1983 Total QE Cross Section

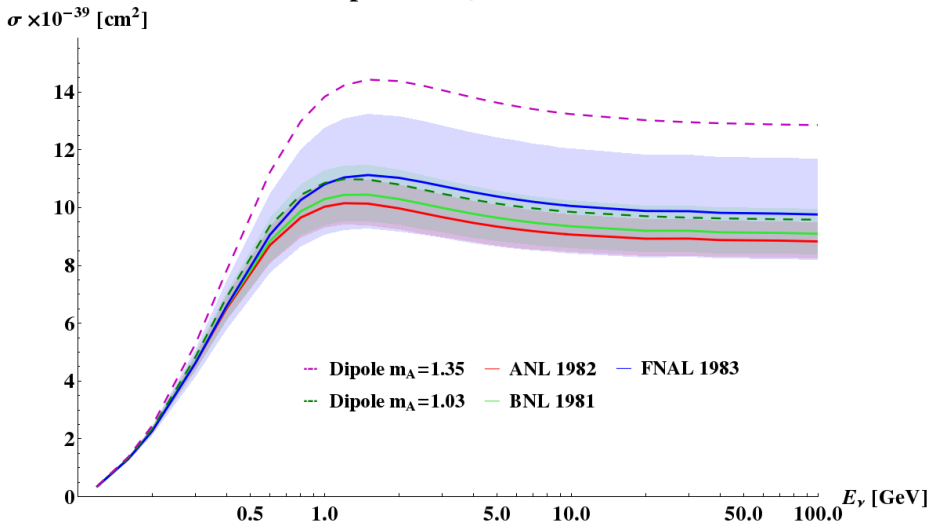
$\sigma \times 10^{-39} [\text{cm}^2]$



N	χ^2/DOF
2	0.78
4	0.85
6	0.95
8	1.09

Deuterium Fitting Results

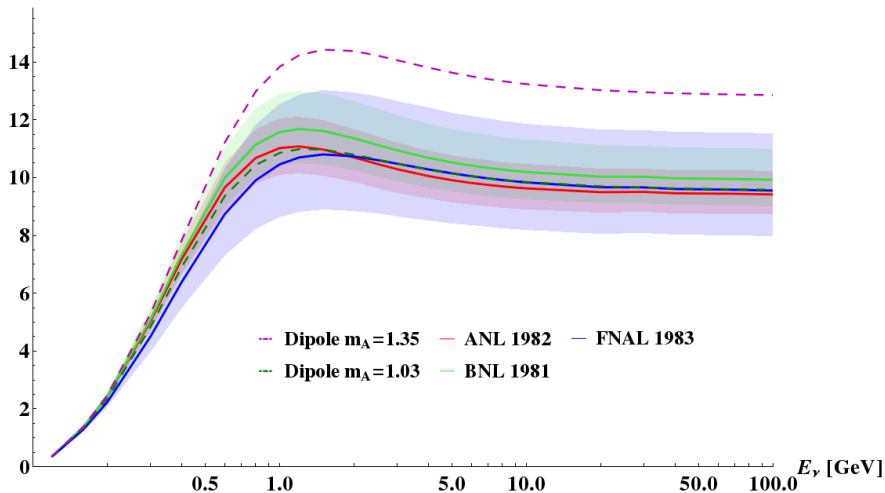
N=2 z-Expansion QE Cross Section



Deuterium Fitting Results

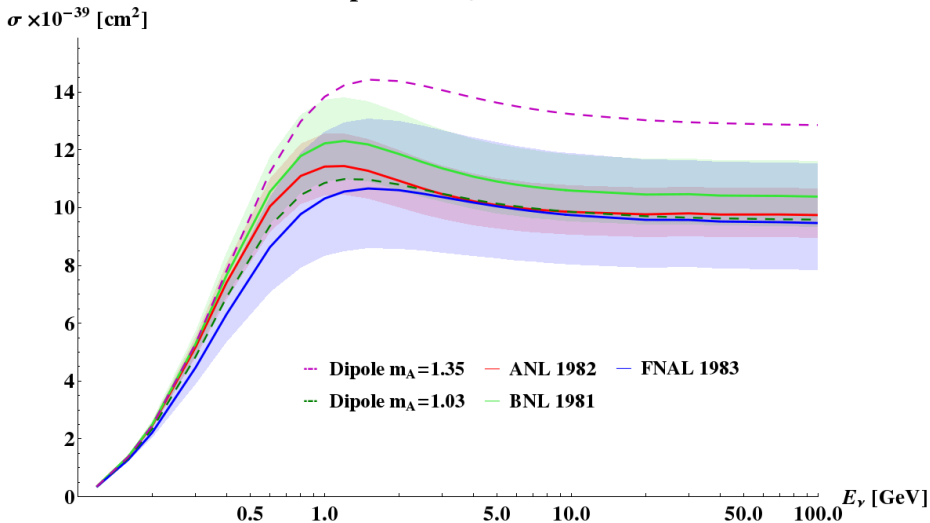
N=4 z-Expansion QE Cross Section

$\sigma \times 10^{-39} [\text{cm}^2]$

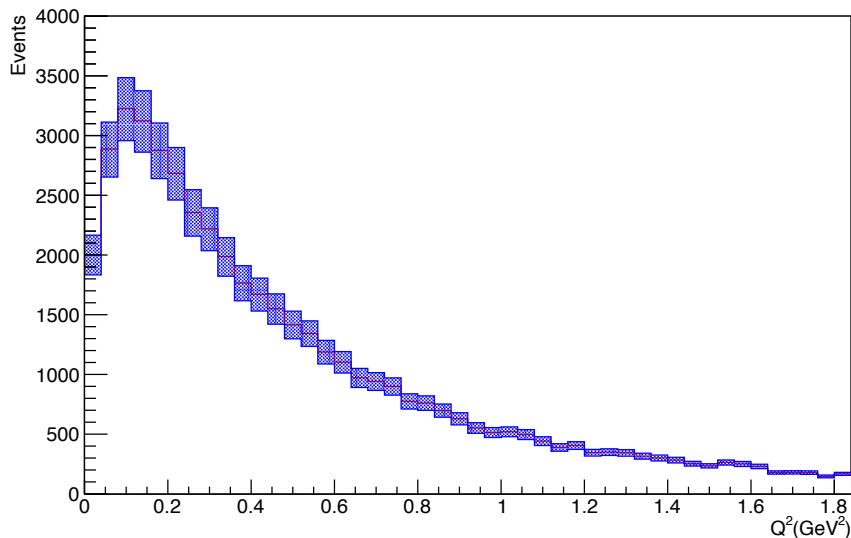


Deuterium Fitting Results

N=6 z-Expansion QE Cross Section



Implications for MINER ν A



Fit to BNL N=3
(Figure made by M. Betancourt)

Take Home Messages

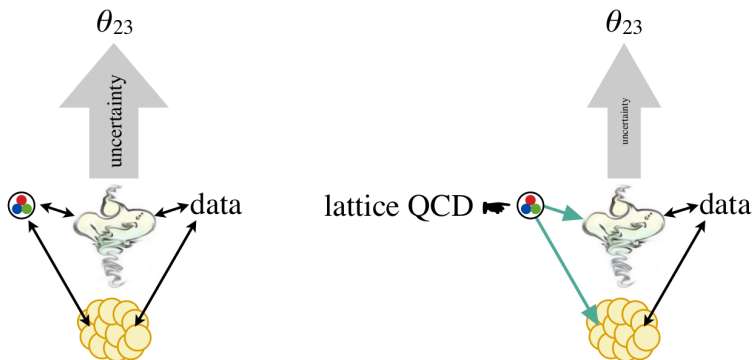
- χ^2/DOF approximately the same for z-expansion and dipole
- FNAL: truncation error negligible for $N \geq 4$
- ANL/BNL: truncation error small by $N = 4$, negligible for $N \geq 6$
- ANL/BNL: Large a_3 , deuterium effects?
- Given choice of priors, errors on total cross section larger than dipole by 1.5-2 times
- Total cross sections from z-expansion fits consistent to within 1σ

Still to come:

- More detailed deuterium corrections
- Radiative corrections
- Study of effects on MINER ν A observables

Lattice QCD in Neutrino Physics

LQCD can play important role in breaking degeneracy



Nucleon-level/Nucleus-level effects entangled

Measurements of observables are **model-dependent**

LQCD acts as **disruptive technology** to break the cycle

Lattice F_A Calculation

$g_A = F_A(q^2 = 0)$ is a historically difficult calculation
(world best 10-20% too small)

What makes it hard:

- Baryons(!)
- Finite size effects
- Chiral extrapolation
- Explicit Chiral symmetry breaking - for some formalisms
- Excited state contamination

Lattice F_A Calculation

Will use MILC's 2+1+1 flavor gauge ensembles

What we bring to the table:

- High computation speed \rightarrow Statistics
- Large lattices \rightarrow Control finite size effects
- Physical quark masses \rightarrow Avoid chiral extrapolation
- Exact chiral symmetry \rightarrow Obtain absolute normalization
- Variational method \rightarrow
Mitigate excited state contamination

Outlook

Deuterium Fitting/GENIE:

- Finish fitting/writeup - next few months
- Write correlated reweighting for GENIE
- Coordinate GENIE code release with publication release

Lattice:

- Code testing/development - this/next month
- Production - soon after
- g_A calculation - Spring/Summer 2016
- $F_A(q^2)$ calculation - Fall/Winter 2016

Further (more challenging) lattice QCD calculations:

- $\nu_e N \rightarrow \nu_e N'$
- N- Δ transition currents
- $\nu_e N \rightarrow \pi \ell N'$
- $\nu_e N \rightarrow \pi \ell \Sigma$

Conclusions

Neutrino physics is subject to

underestimated and model-dependent systematics

- To reduce **systematics from modeling**, need to understand **nuclear physics**
- To understand **nuclear physics**, need to understand **nucleon-level cross sections** from an ab initio calculation
 - z-Expansion removes model assumptions and permits better understanding of systematic errors
 - hydrogen (deuterium) targets have [almost] no nuclear effects
 - LQCD offers a way to access nucleon form factors directly

Thanks!

Backup Slide(s)

Error Budgets

Source of Uncertainty	MINOS Absolute/ ν_e	T2K ν_e	LBNE ν_e	Comments
Beam Flux after N/F extrapolation	3%/0.3%	2.9%	2%	MINOS is normalization only. LBNE normalization and shape highly correlated between ν_μ/ν_e .
Detector effects				
Energy scale (ν_μ)	7%/3.5%	included above	(2%)	Included in LBNE ν_μ sample uncertainty only in three-flavor fit. MINOS dominated by hadronic scale.
Absolute energy scale (ν_e)	5.7%/2.7%	3.4% includes all FD effects	2%	Totally active LArTPC with calibration and test beam data lowers uncertainty.
Fiducial volume	2.4%/2.4%	1%	1%	Larger detectors = smaller uncertainty.
Neutrino interaction modeling				
Simulation includes: hadronization cross sections nuclear models	2.7%/2.7%	7.5%	~ 2%	Hadronization models are better constrained in the LBNE LArTPC. N/F cancellation larger in MINOS/LBNE. X-section uncertainties larger at T2K energies. Spectral analysis in LBNE provides extra constraint.
Total	5.7%	8.8%	3.6 %	Uncorrelated ν_e uncertainty in full LBNE three-flavor fit = 1-2%.

LBNE Experiment