

Self-Consistent Treatment of Neutrino Physics in Cosmology

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17 Jul 2015

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UCSD

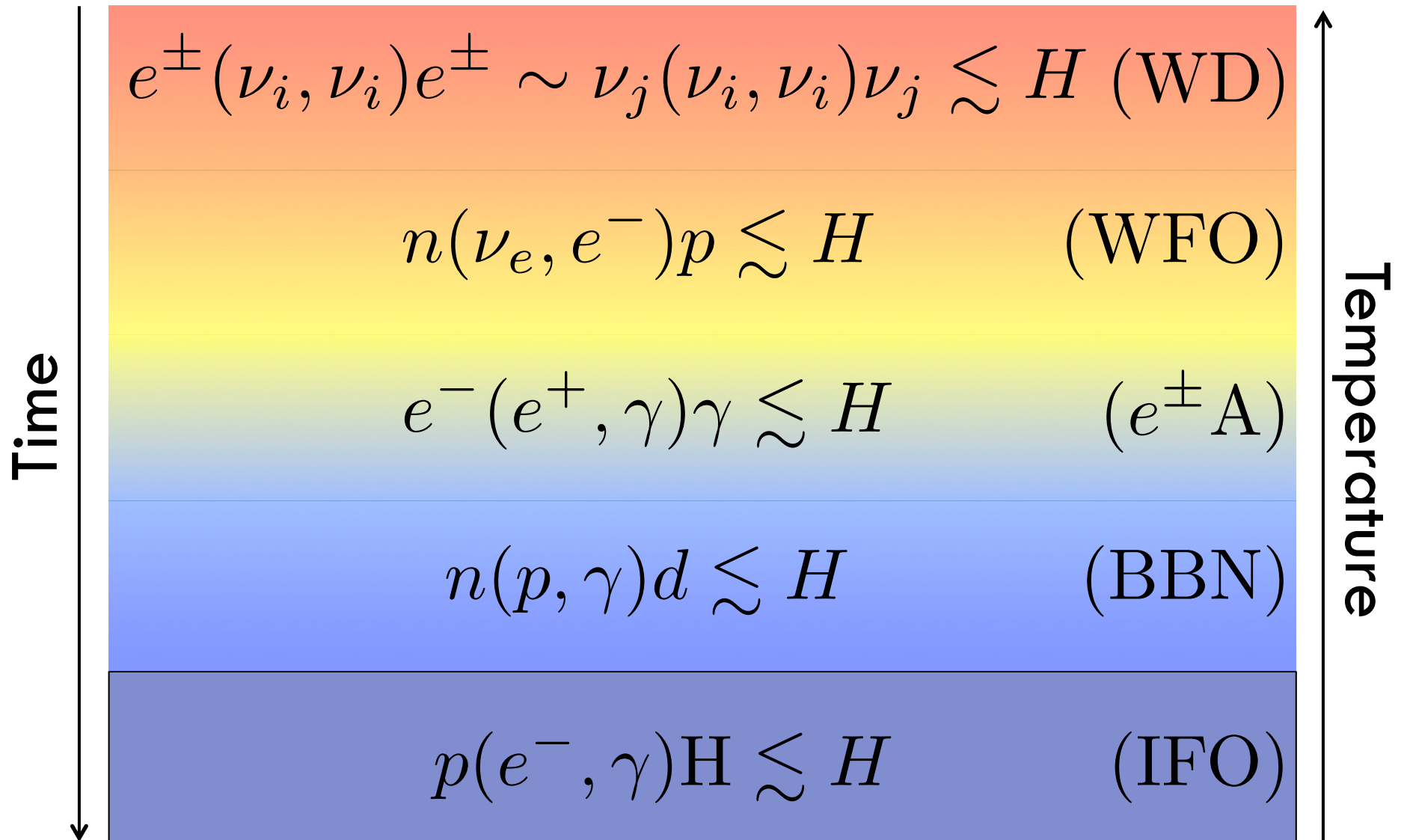


Outline



- Overview of the early universe
- An alternative to N_{eff}
- The “Neutrino-Mass/Recombination Effect” (VMR)
- Neutrino Transport during Weak Decoupling and BBN

Epochs of Interest



Code

B

BN

Predict primordial nuclear abundances

U

NITARY

Preserve unitarity in nuclear reaction network

Quantify errors

R

ECOMBINATION

Treat recombination with three-level atom similar to recfast [16]

Isolate neutrino signatures in cosmological power spectra

S

ELF-CONSISTENT

Maintain self-consistency over large range of epochs

T

RANSPORT

Follow evolution of neutrino spectra



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Neutrino Energy Density

$$\rho_{\text{rad}} = \left(2 + \frac{7}{4} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right) \frac{\pi^2}{30} T_{\gamma}^4$$
$$\implies N_{\text{eff}} = 3$$

Finite temperature radiative correction

Dicus et al. (1982) [6]

Lopez & Turner (1999)[7]

$$\Delta N_{\text{eff}} \sim 0.011$$

Weak decoupling is not a sharp event

Dolgov, Hansen, & Semikoz (1997)[8]

$$\Delta N_{\text{eff}} \sim 0.034$$

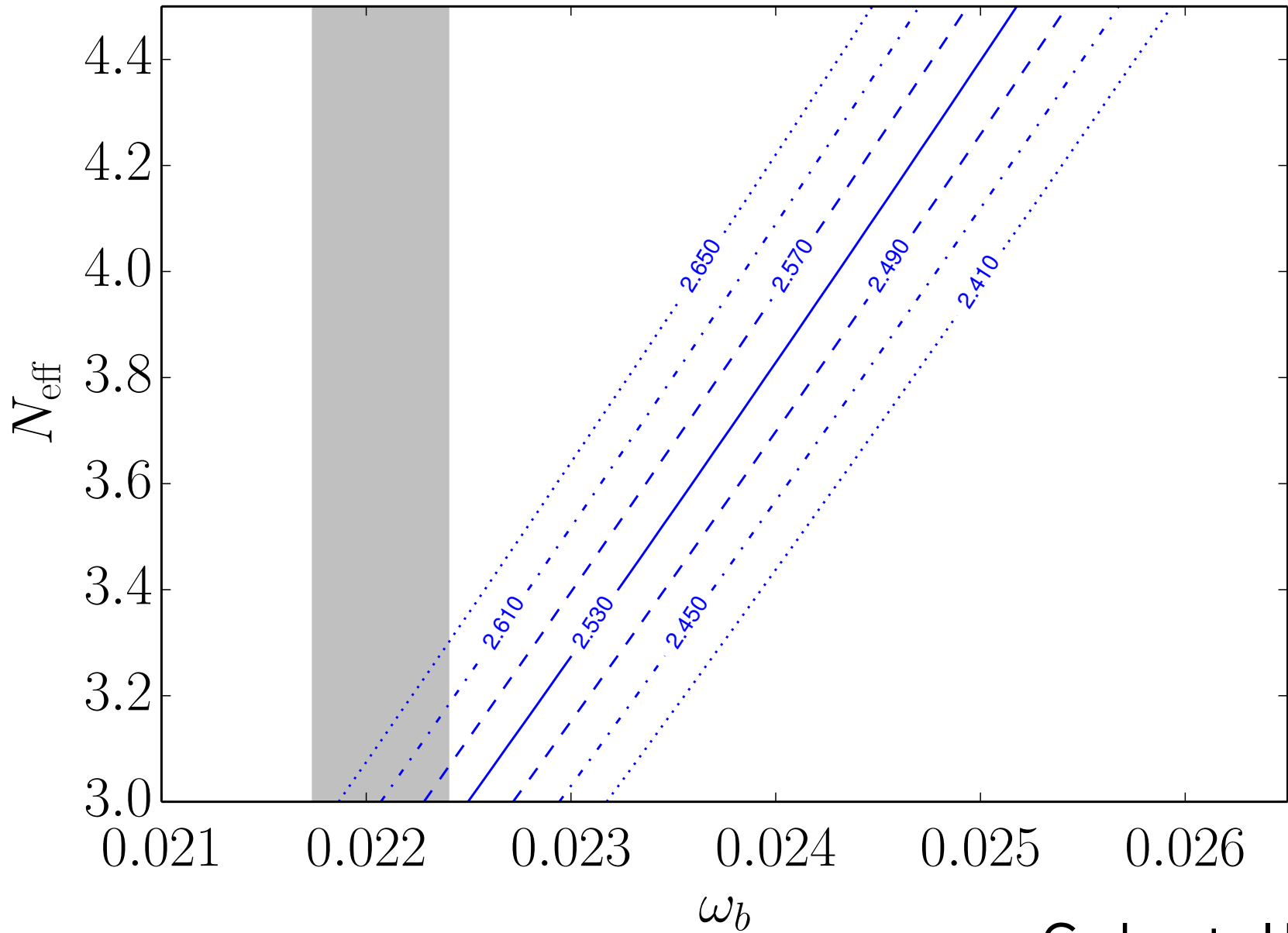
Mangano et al. (2005)[9]:

$$N_{\text{eff}} = 3.046 \text{ in SM}$$

Planck XIII (2015)[10]:

$$N_{\text{eff}} = 3.15 \pm 0.23$$

Contours of constant D/H



Grohs, et al.[13]

Issues with N_{eff}

- N_{eff} not measured immediately after e^\pm annihilation or during BBN
- ρ_{rad} is not a CMB observable; how does CMB determine ρ_{rad} ?
- Neutrinos have non-zero rest mass.
- What happens to N_{eff} when including Beyond-Standard-Model physics?

$$N_{\text{eff}}^{(\text{th})} \text{ vs. } \tilde{N}_{\text{eff}} : \begin{cases} \rho_\gamma + \rho_\nu = \left(1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}}^{(\text{th})} \right) \rho_\gamma \\ \tilde{N}_{\text{eff}} = N_{\text{eff}}^{(\text{th})} [r_s/r_d = (r_s/r_d)^{(\text{inp})}] \end{cases}$$

Strategy

Sound Horizon:

$$r_s = \int_0^{a_{\gamma d}} \frac{da}{a^2 H \sqrt{3(1+R)}} \quad R \equiv \frac{3\rho_b}{4\rho_\gamma}$$

Damping Diffusion Wave Number:

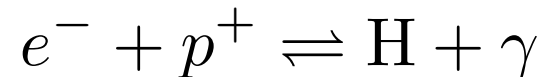
$$k_d^{-2} = \int_0^{a_{\gamma d}} \frac{da}{n_e \sigma_T H a^3 6(1+R)} \left(\frac{R^2}{1+R} + \frac{16}{15} \right), \quad r_d = \pi/k_d$$

$$H = \left(\frac{8\pi}{3m_{\text{pl}}^2} \rho \right)^{1/2} \rightarrow \left(\frac{8\pi}{3m_{\text{pl}}^2} (\rho_m + \rho_\Lambda + \rho_{\text{rad}}) \right)^{1/2}$$

Ratio: $r_s/r_d = d_s/d_d = \theta_s/\theta_d, \quad \theta \equiv \frac{d}{D_A}$

Free-Electron Fraction

Typical reaction of Interest:



Covariant form of Boltzmann Eqn.:

$$p^{\sigma} \frac{\partial f}{\partial x^{\sigma}} - \Gamma_{\mu\nu}^{\sigma} p^{\mu} p^{\nu} \frac{\partial f}{\partial p^{\sigma}} = \hat{C}[f]$$

Evolution of X_e :

$$\frac{dX_e}{dt} = (1 - X_e)\beta - X_e^2 n_b \alpha^{(2)}$$

β : Ionization coefficient

n_b : Baryon number density

$\alpha^{(2)}$: Recombination coefficient

Boltzmann Eqn. Correction

Correction to Boltzmann Equation due to Peebles (1968):

$$\frac{dX_e}{dt} = \left[(1 - X_e)\beta - X_e^2 n_b \alpha^{(2)} \right] C$$

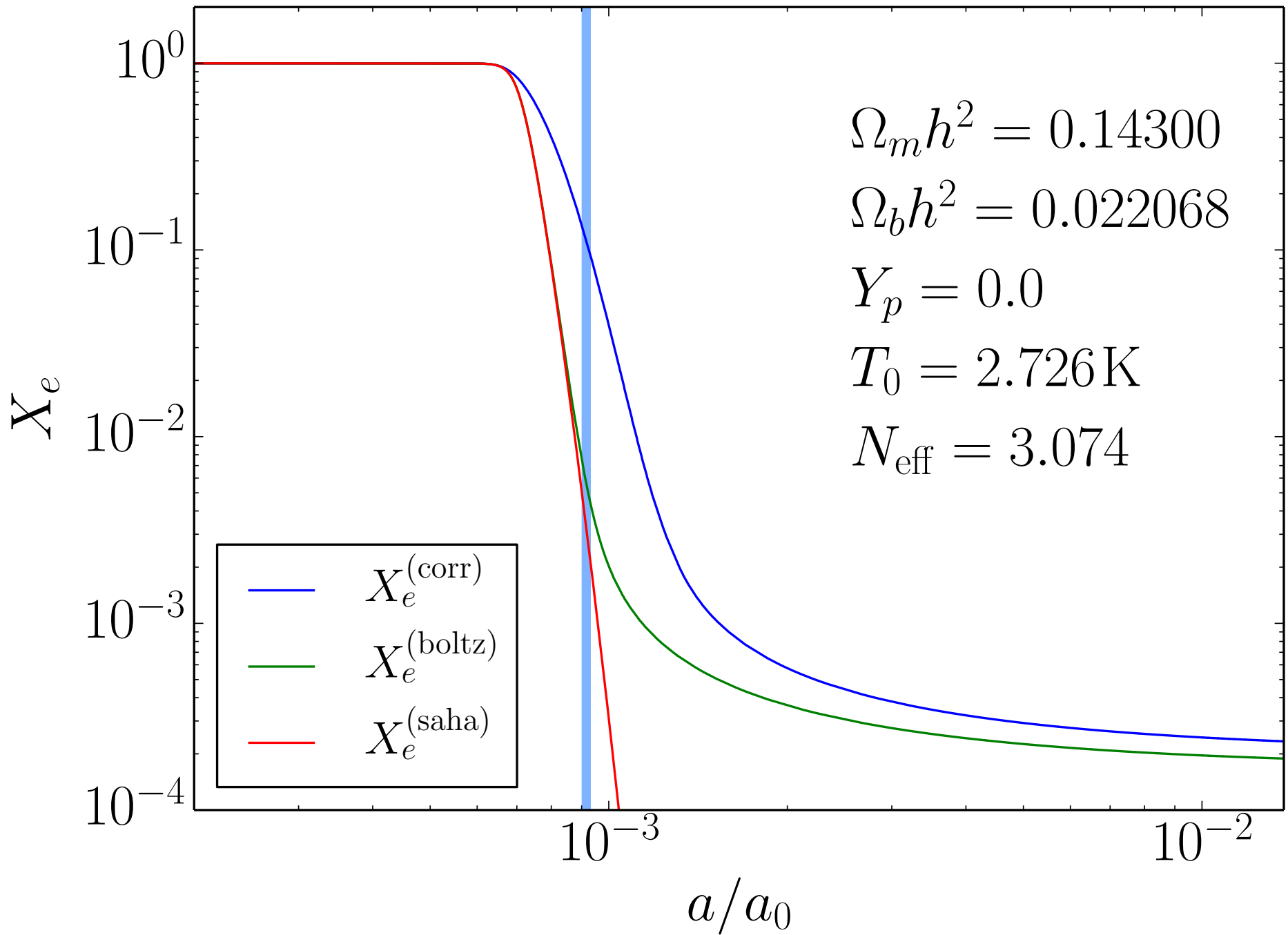
Correction Factor for 3-level atom:

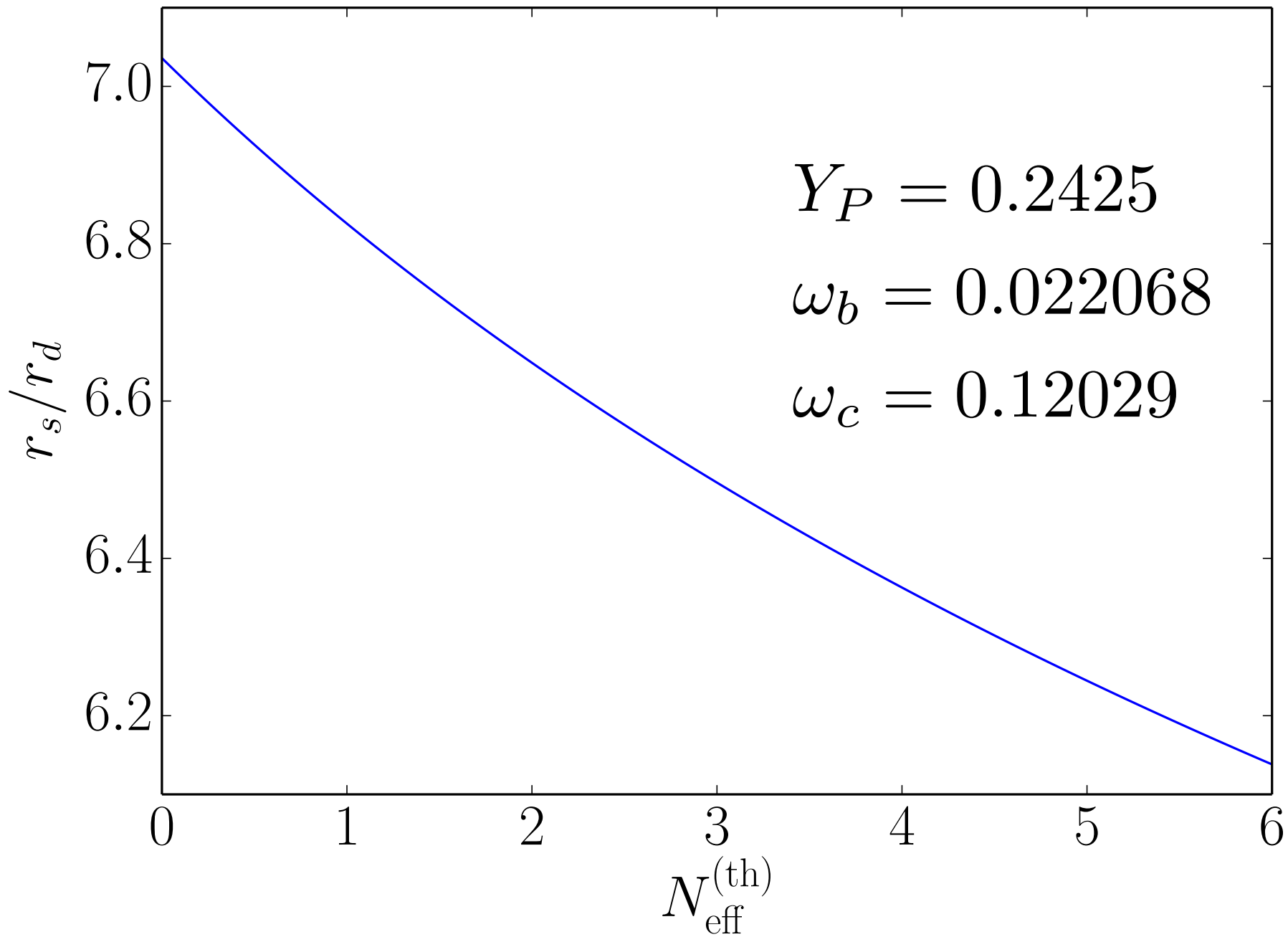
$$C = \frac{\Lambda_\alpha + \Lambda_{2\gamma}}{\Lambda_\alpha + \Lambda_{2\gamma} + \beta^{(2)}}$$

Λ_α : Lyman-alpha photon redshift rate

$\Lambda_{2\gamma}$: 2-photon decay rate

$\beta^{(2)}$: Lyman-alpha production rate





ν MR effect

Neutrinos free-stream after weak decoupling:

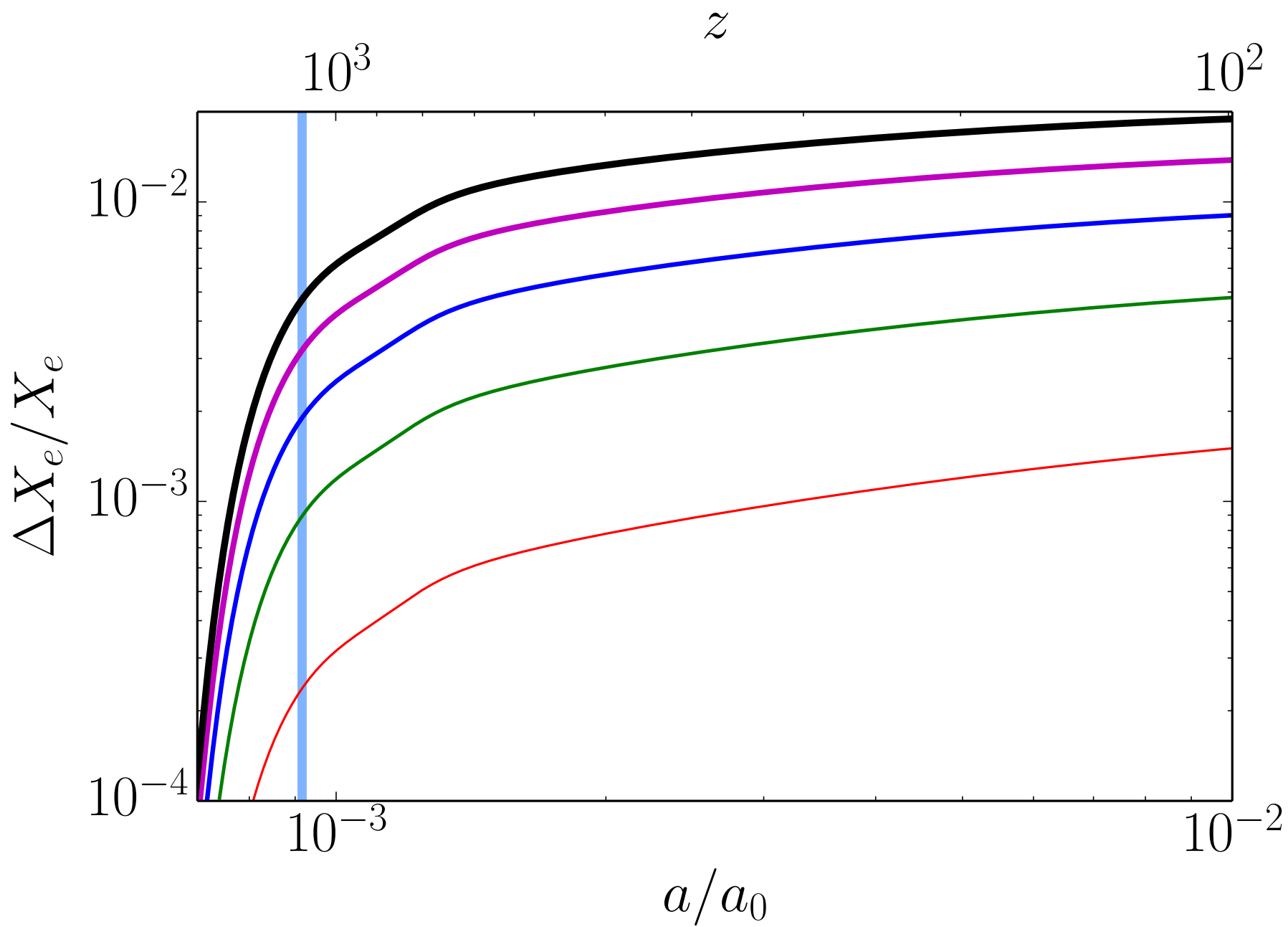
$$\rho_\nu(m \neq 0) = \int_0^\infty \frac{p^2 \sqrt{p^2 + m^2} dp}{e^{p/T} + 1} > \int_0^\infty \frac{p^3 dp}{e^{p/T} + 1} = \rho_\nu(m = 0)$$

Naïve Scaling Argument:

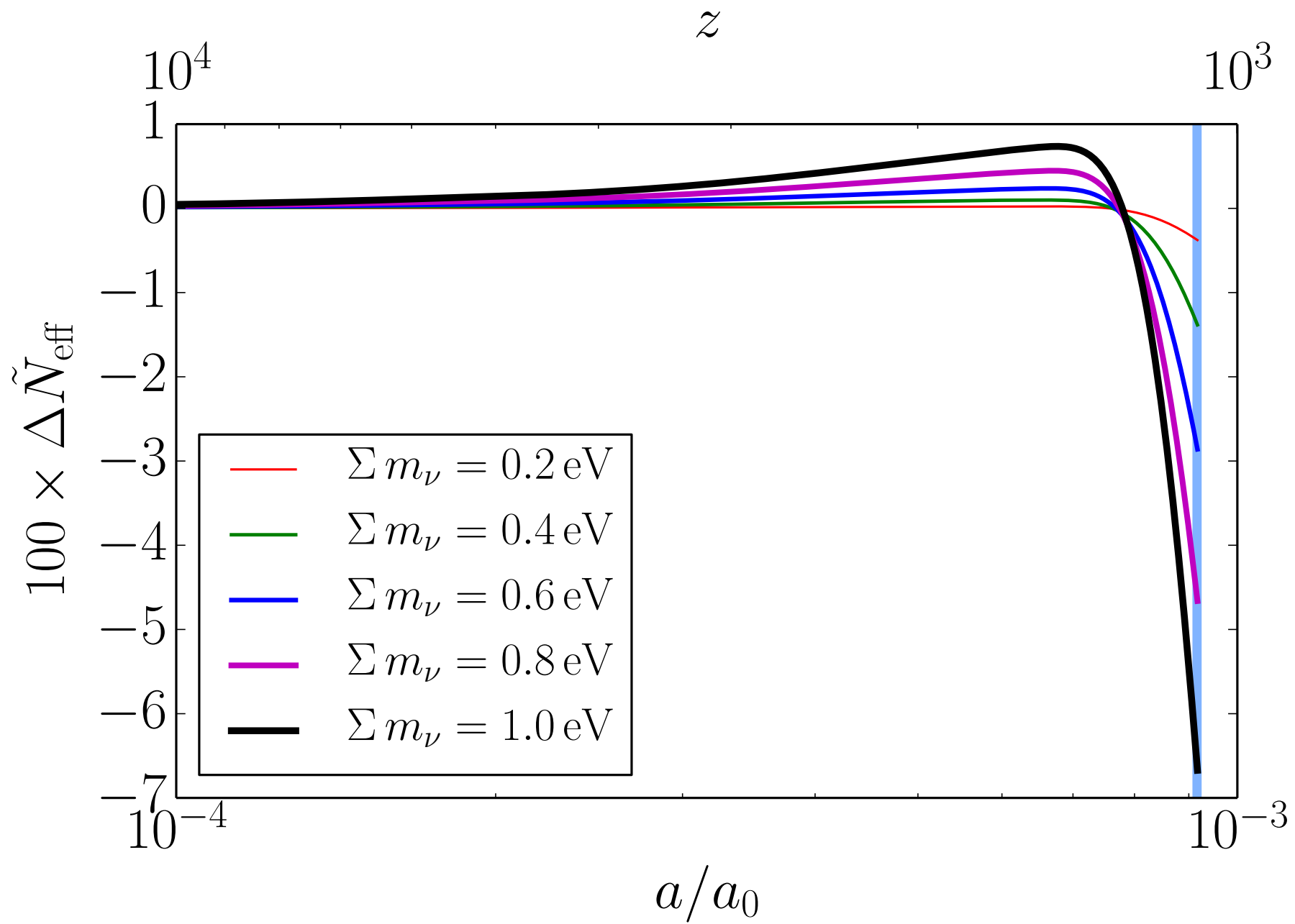
$$r_s \propto \frac{1}{H}, r_d \propto \frac{1}{\sqrt{H}} \implies \frac{r_s}{r_d} \propto \frac{1}{\sqrt{H}} \sim \frac{1}{\rho^{1/4}}$$

Expectation: Larger Hubble rate implies Larger \tilde{N}_{eff}

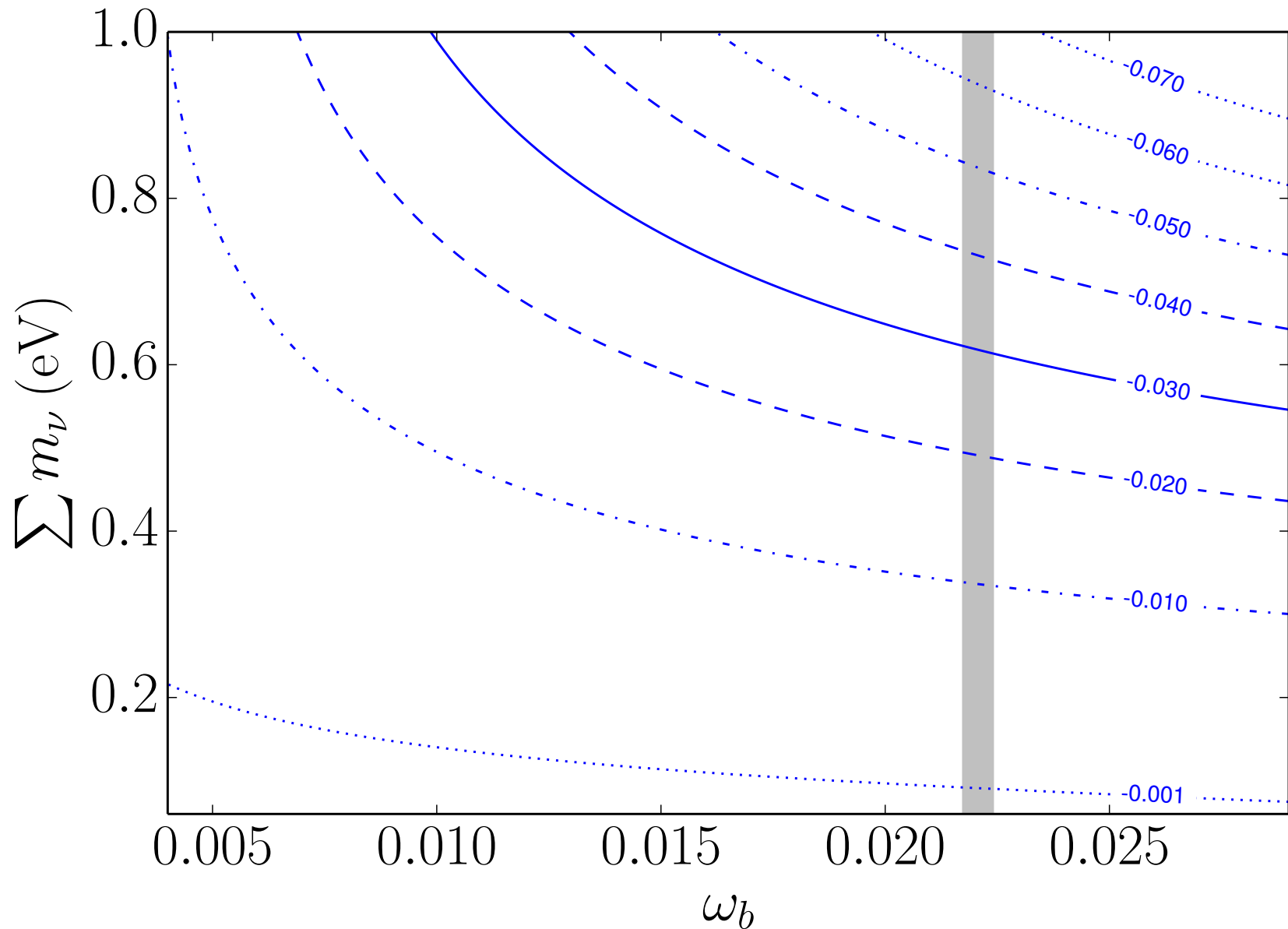
However, calculation yields: $\tilde{N}_{\text{eff}} = 2.995$, $\Sigma m_\nu = 0.23 \text{ eV}$



Grohs, et al.[12]



Contours of constant $\Delta\tilde{N}_{\text{eff}}$

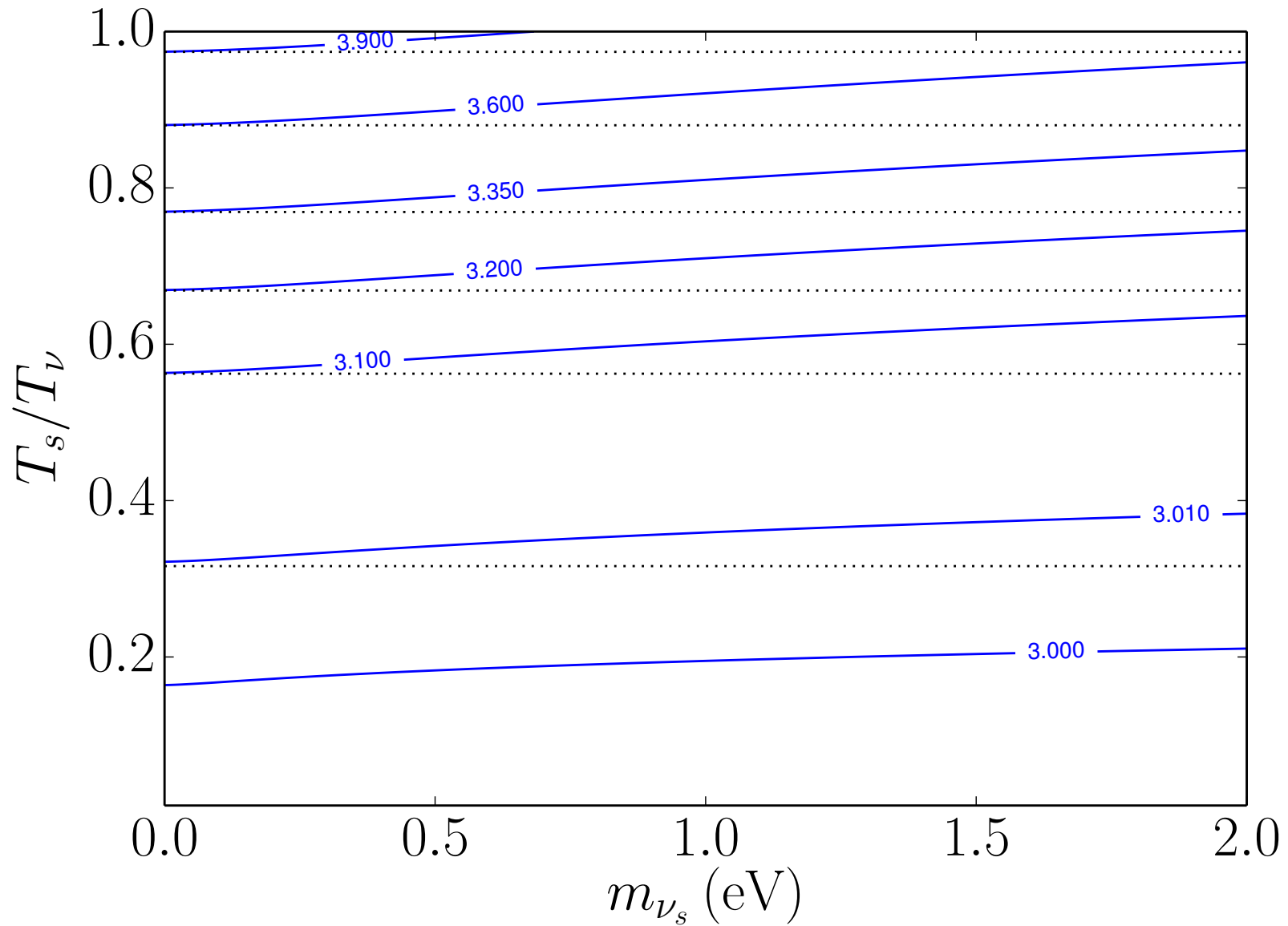


Grohs, et al.[13]

Hints for light sterile neutrinos?

- mini-BooNE
 - neutrino oscillation experiment $\nu_e \rightarrow \nu_s \rightarrow \nu_\mu$
 - appearance with $\delta m^2 \sim 1 \text{ eV}^2$
 - result inconsistent with flavor oscillation alone
- Neutrino reactor anomaly
 - 3σ deficit neutrinos detected in short-baseline (<100m) reactor experiments
 - $\bar{\nu}_e$ deficit from $\bar{\nu}_e \rightarrow \bar{\nu}_s$ (???) – a disappearance experiment
 - A. Hayes et al. (2013)[15] find “large corrections”
- Extra radiation at photon-decoupling (Neff) ??
 - CMB observations (PolarBear, ACT, SPT, Planck, CMBPol,...)
 - ‘extra’ RED could reconcile H_0 and σ_8 inferred from CMB and astronomical observation

Contours of constant \tilde{N}_{eff}



Weak Interactions

Channels:

$$\nu_i + \nu_j \leftrightarrow \nu_i + \nu_j$$

$$\nu_i + \bar{\nu}_j \leftrightarrow \nu_i + \bar{\nu}_j$$

$$\nu_i + \bar{\nu}_i \leftrightarrow \nu_j + \bar{\nu}_j$$

$$\nu_i + e^\pm \leftrightarrow \nu_i + e^\pm$$

$$\bar{\nu}_i + e^\pm \leftrightarrow \bar{\nu}_i + e^\pm$$

$$\nu_i + \bar{\nu}_i \leftrightarrow e^- + e^+$$

Summed-Squared Amplitude examples:

$$\nu_e(1) + \nu_e(2) \leftrightarrow \nu_e(3) + \nu_e(4)$$

$$\langle |\mathcal{M}|^2 \rangle = 2^7 G_F^2 (P_1 \cdot P_2)^2$$

$$\nu_e(1) + e^-(2) \leftrightarrow e^-(3) + \nu_e(4)$$

$$\begin{aligned} \langle |\mathcal{M}|^2 \rangle = 2^5 G_F^2 & [(1 + 2 \sin^2 \theta_W)^2 (P_1 \cdot Q_2)(Q_3 \cdot P_4) \\ & + 4 \sin^4 \theta_W (P_1 \cdot Q_3)(Q_2 \cdot P_4) \\ & - 2 \sin^2 \theta_W (1 + 2 \sin^2 \theta_W) m_e^2 (P_1 \cdot P_4)] \end{aligned}$$

Boltzmann Equation Revisited

Covariant form of Boltzmann Eqn.:

$$p^\sigma \frac{\partial f}{\partial x^\sigma} - \Gamma_{\mu\nu}^\sigma p^\mu p^\nu \frac{\partial f}{\partial p^\sigma} = \hat{C}[f]$$

In comoving coordinates: $f(\vec{x}, \vec{p}, t) \rightarrow f(p, t) \rightarrow f(\epsilon, t)$

With collision operator:

$$\begin{aligned} \frac{Df_1}{Dt} = & \int \frac{s}{2E_1} \frac{d^3q_2}{(2\pi)^3 2E_2} \frac{d^3q_3}{(2\pi)^3 2E_3} \frac{d^3p_4}{(2\pi)^3 2E_4} \\ & \times \langle |\mathcal{M}|^2 \rangle (2\pi)^4 \delta^4(P_1 + Q_2 - Q_3 - P_4) \\ & \times [f_3 f_4 (1 - f_1)(1 - f_2) - f_1 f_2 (1 - f_3)(1 - f_4)] \end{aligned}$$

Collision Term Reduction

1. Nine-dimensional integral over phase space of particles 2, 3, and 4
2. Conservation of four-momentum – Five-dimensional integral
3. Isotropy – Three-dimensional integral
4. Integration Limits Trick – Two-dimensional integral

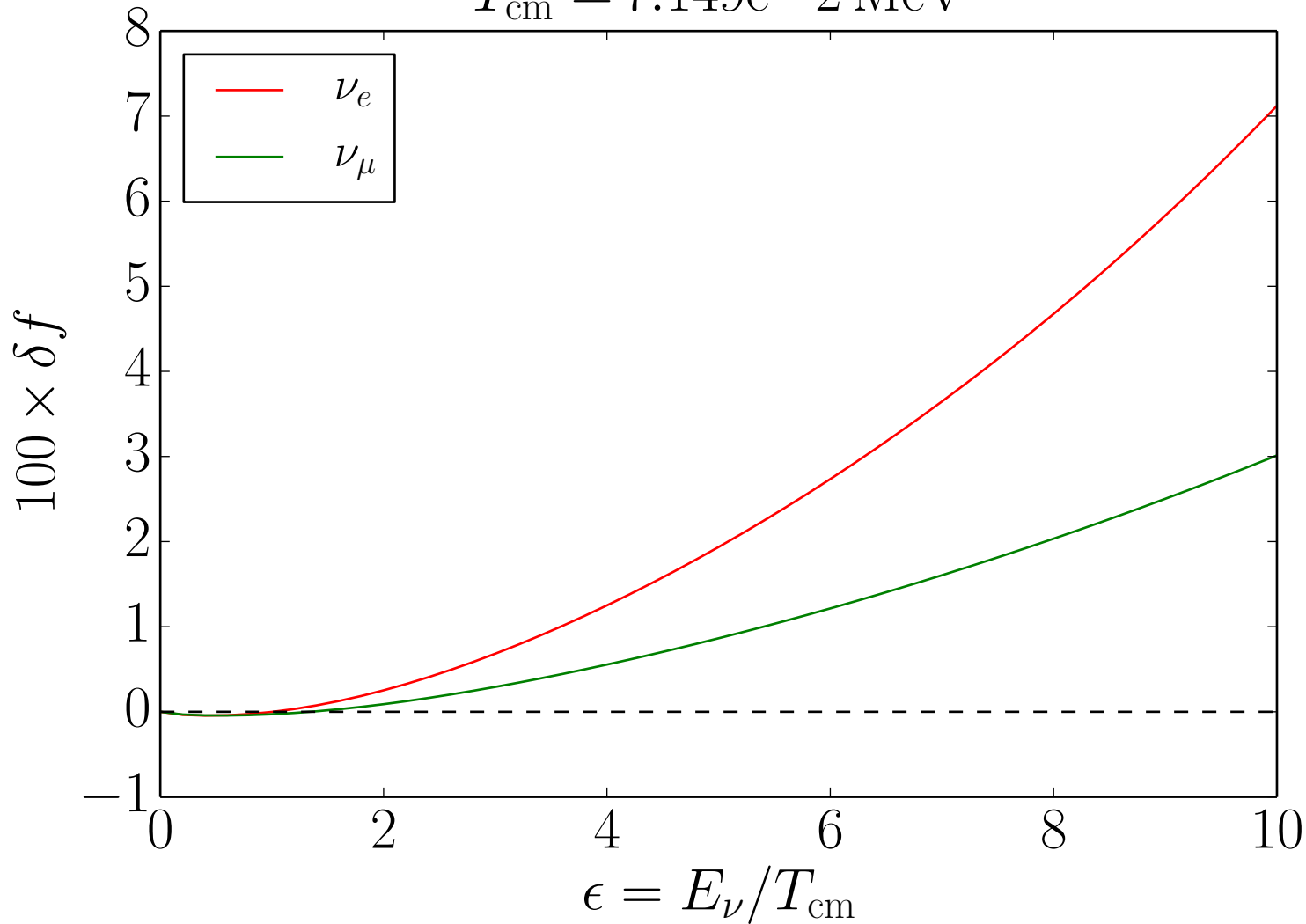
$$\frac{Df_1}{Dt} = \frac{\kappa}{32(2\pi)^3} \int_0^\infty dp_2 p_1 p_2^3 \int_0^{p_1+p_2} dp_3 W(p_1, p_2, p_3) F(p_1, p_2, p_3, p_1 + p_2 - p_3)$$
$$W(p_1, p_2, p_3) = \int_{x_0}^1 dx \frac{(1-x)^2}{\sqrt{p_1^2 + p_2^2 + 2p_1 p_2 x}}$$
$$x_0 = \max\left(-1, 1 - \frac{2p_3(p_1 + p_2 - p_3)}{p_1 p_2}\right)$$

Computational Parameters

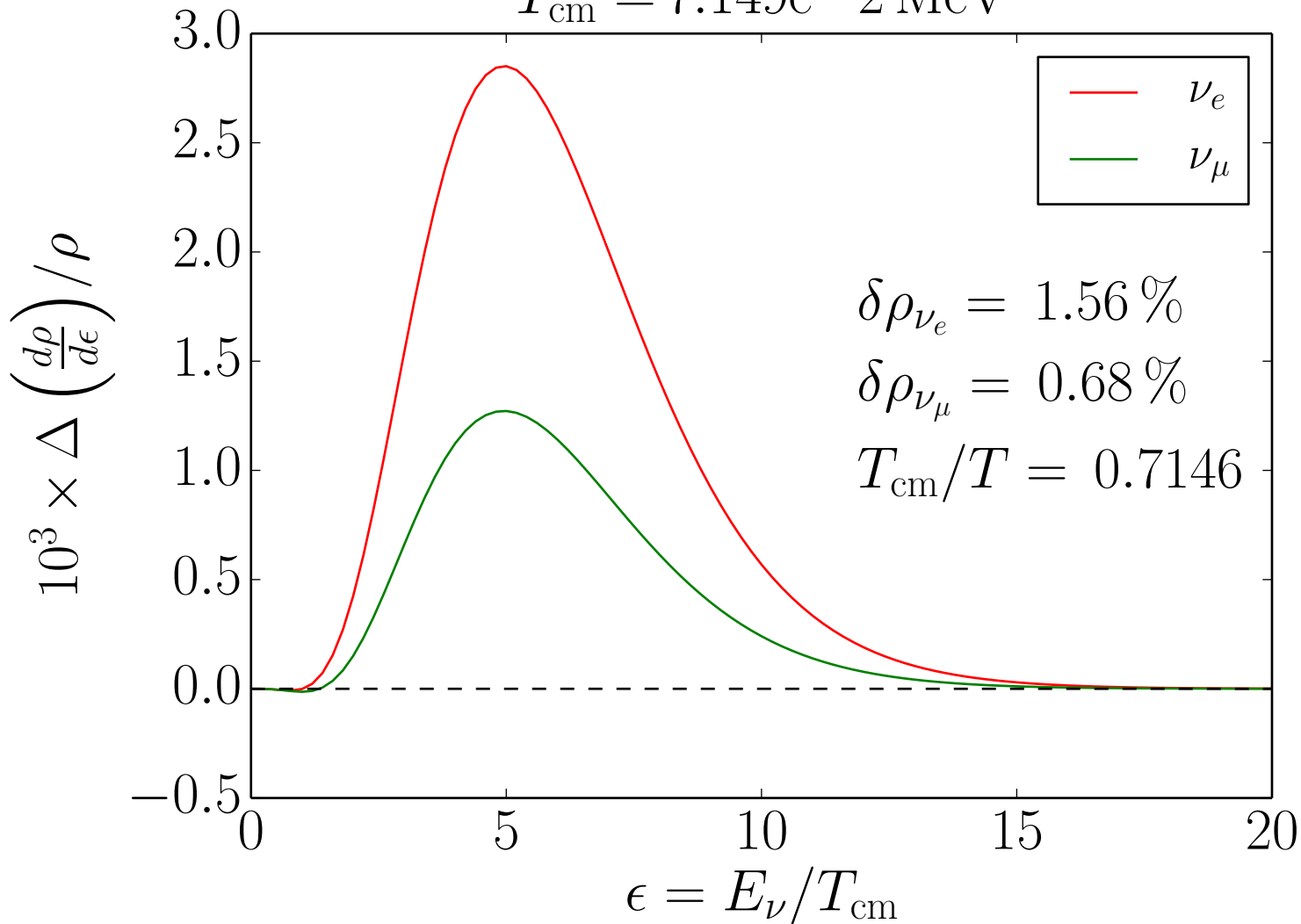


1. All calculations done in the weak eigenbasis.
2. 100 bins, 101 abscissas linearly spaced between $0 \leq \epsilon \leq 20$.
3. Input temperature $T_{\text{in}} = 8 \text{ MeV}$, final temperature $T_{\text{cm}} \sim 9 \text{ keV}$.
4. Fifth-order polynomial interpolator.
5. Acceptance tolerance $(\text{net/frs})^{(\text{tol})} = 30$.
6. Electrons and positrons always in thermal/chemical equilibrium.

$$T_{\text{cm}} = 7.149\text{e-}2 \text{ MeV}$$



$$T_{\text{cm}} = 7.149\text{e}-2 \text{ MeV}$$



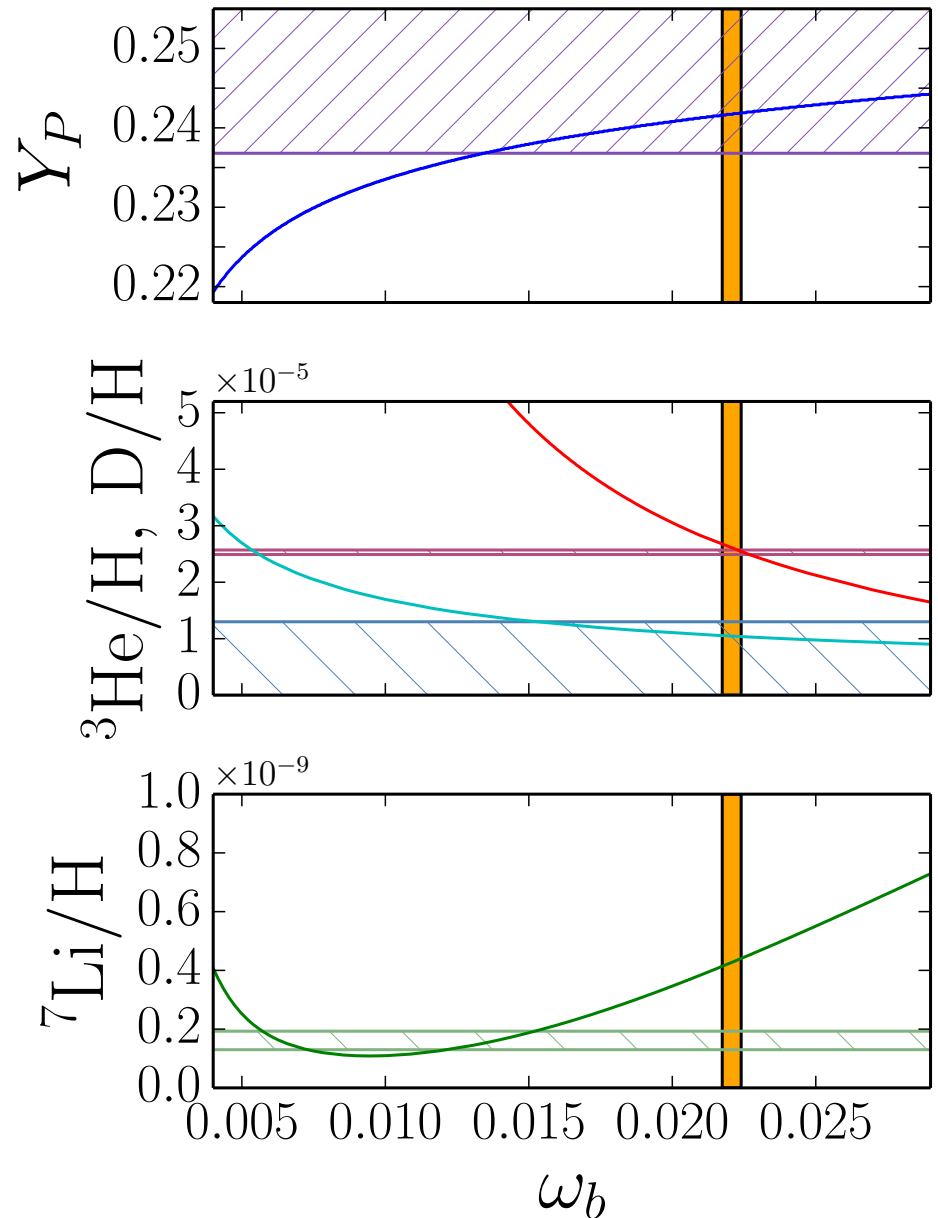
Nucleosynthesis

Without neutrino transport
(sharp WD):

$$Y_P = 0.2425,$$
$$D/H = 2.615 \times 10^{-5},$$

With neutrino transport:

$$Y_P = 0.2426, \delta(Y_P) \simeq 4 \times 10^{-4}$$
$$D/H = 2.628 \times 10^{-5},$$
$$\implies \delta(D/H) \simeq 5 \times 10^{-3}$$



Summary and Future Work



→ BURST

→ Effect of neutrino rest mass on ionization equilibrium freeze-out.
arXiv: 1412.6875

→ Probing neutrino physics with a self-consistent treatment of the weak decoupling, nucleosynthesis, and photon decoupling epochs.
arXiv: 1502.02718

→ Beyond Standard Model Physics

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