

This lecture lies a bit outside the line of our course, in that we shall consider a non-mean field setting. More precisely, we let $G = (V, E)$ be a square grid of side \sqrt{N} and consider the Ising model on such a grid:

$$\mu(x) = \frac{1}{Z(\beta)} \exp \left\{ \beta \sum_{(ij) \in E} x_i x_j \right\}, \quad (1)$$

and will prove the following result.

Theorem 1. *Let X be a random configuration from the two-dimensional Ising model (1) and $\bar{X} = N^{-1} \sum_{i=1}^N X_i$ denote its magnetization.*

Then there exist two inverse temperatures $\beta_1 > \beta_2 > 0$, such that the following happens. If $\beta > \beta_1$ ('low-temperature'), then there exist $\delta(\beta) > 0$ such that

$$\mathbb{E}|\bar{X}| \geq \delta(\beta). \quad (2)$$

If $\beta < \beta_2$ ('high-temperature'), then, for any $\delta > 0$

$$\lim_{N \rightarrow \infty} \mathbb{P}\{|\bar{X}| \leq \delta\} = 1. \quad (3)$$

Our proofs will be based on two classical arguments as found in [Gri64] (low temperature) and [Fis67] (high temperature). While the above statement and our proofs refer to the two-dimensional case, the techniques are more general.

1 Low temperature: Peierls argument

The *dual lattice* of G is the graph $G^* = (V^*, E^*)$ constructed as follows. Imagine to draw G on a plane, using the standard embedding of \mathbb{Z}^2 in \mathbb{R}^2 . Then, for each edge (i, j) in the original lattice G , draw a perpendicular edge of length 1, centered at the midpoint of (i, j) . These are the edges of the G^* . Its vertices are the endpoints of such edges (only whenever two endpoints coincide, the vertices do coincide as well).

The number of vertices in the dual lattice is $|V^*| = (\sqrt{N} + 1)^2 - 4$.

A *Peierls contours* configuration is a set $\mathfrak{C} = \{C_1, \dots, C_n\}$ of directed paths C_i in G^* that satisfy the following conditions:

- (i) Each path is either closed or has on the boundary (degree-one vertices) of G_* .
- (ii) No two paths cross each other.
- (iii) Whenever two paths share a vertex, each of them bends to the right.
- (iv) Whenever path C_i is 'directly inside' path C_j , the two paths have opposite orientations.

In the last point, by 'directly inside' we mean the following. If C_i and C_j are closed paths, then it must be that C_i is inside C_j but not inside any other path C_k which in turn is inside C_j . If they are not closed, we first close them around the lattice and then require the same condition to be satisfied.

To a given Peierls contours configuration we associate the probability

$$\mu_*(\mathfrak{C}) = \frac{1}{Z^*(\beta)} \prod_{C \in \mathfrak{C}} e^{-2\beta|C|}. \quad (4)$$

Lemma 2. *There is a one-to-one correspondence between Ising model configurations and Peierls contours configurations. Further, under this correspondence, the Ising measure $\mu(x)$ is mapped to the Peierls one $\mu_*(\mathfrak{C})$.*

Given an Ising configuration x , the corresponding Peierls contours separate the vertices with $x_i = +1$ from those with $x_i = -1$.

Proof By figures. □

Proof [Theorem 1, low temperature.] For the sake of simplicity we shall prove that $\mathbb{P}\{|\bar{X}| = 0\}$ is bounded as in the statement. The cases $\mathbb{P}\{|\bar{X}| = k\}$ with $|k| \leq N\delta$ can be treated analogously and the thesis (2) follows by union bound.

First notice that the number of Peierls contours C with $|C| = l$ is upper bounded as

$$n(l) \leq |V^*| 3^l \leq N 3^{l+1}. \quad (5)$$

This is showed by first choosing a starting point (for which there are $|V^*|$ possibilities), and then a length- l non-reversing walk in \mathbb{Z}^2 .

The second remark is that, for any contour C

$$\mu_*(C \in \mathfrak{C}) \leq e^{-2\beta|C|}, \quad (6)$$

where $|C|$ denotes the length of C . This is proved by writing

$$\mu_*(C \in \mathfrak{C}) = \frac{\sum_{\mathfrak{C} \ni C} e^{-2\beta|\mathfrak{C}|}}{\sum_{\mathfrak{C}'} e^{-2\beta|\mathfrak{C}'|}}, \quad (7)$$

and noticing that, for each configuration \mathfrak{C} in the numerator, there exists a configuration $\mathfrak{C}' = \mathfrak{C} \setminus C$ in the denominator. The bound follows by restricting the sum in the denominator to such terms and using $|\mathfrak{C}'| = |\mathfrak{C}| - |C|$.

Given a contours configuration \mathfrak{C} , any contour $C \in \mathfrak{C}$ divides V in two subsets: those *inside* C and those *outside* C . If C is closed, the definition of ‘inside’ and ‘outside’ is obvious. If it is not closed, we will call ‘inside’ the smallest of the two subsets in which C divides V (an arbitrary convention can be chosen if they are equal).

We now divide the set of Ising configurations in two subsets

P(+): All the vertices $i \in V$ with $x_i = -1$ lie inside some contour.

P(-): All the vertices $i \in V$ with $x_i = +1$ lie inside some contour.

It is clear that P(+): and P(-): are in one-to one correspondence under the mapping $\{x_i\} \mapsto \{-x_i\}$. Further any configuration x is either in P(+): or in P(-):.

If we denote by $\mathbb{I}_{P(\pm)}$ the indicator function on P(\pm), and by N_- the number of vertices $i \in V$, with $x_i = -1$, we obtain

$$\mathbb{E}|\bar{X}| = \mathbb{E}\{\bar{X} | \mathbb{I}_{P(+)}\} + \mathbb{E}\{\bar{X} | \mathbb{I}_{P(-)}\} \geq \mathbb{E}\{\bar{X} | \mathbb{I}_{P(+)}\} - \mathbb{E}\{\bar{X} | \mathbb{I}_{P(-)}\} = \quad (8)$$

$$= 2\mathbb{E}\{\bar{X} | \mathbb{I}_{P(+)}\} = 1 - \frac{2}{N}\mathbb{E}\{N_- | \mathbb{I}_{P(+)}\}. \quad (9)$$

If $x \in P(+)$, N_- is upper bounded by the number of vertices that are inside at least one contour, which is in turn not larger than $\sum_{C \in \mathfrak{C}} |C_i|^2$. Therefore

$$\begin{aligned} \mathbb{E}|\bar{X}| &\geq 1 - \frac{2}{N}\mathbb{E}\left\{\sum_{C \in \mathfrak{C}} |C|^2\right\} = 1 - \frac{2}{N}\sum_C |C|^2 \mu_*(C \in \mathfrak{C}) \geq \\ &\geq 1 - \frac{2}{N}\sum_{\ell \geq 2} \ell^2 n(\ell) e^{-2\beta\ell} \geq 1 - 2\sum_{\ell \geq 2} \ell^2 (3e^{-2\beta})^\ell. \end{aligned}$$

It is clear that the last expression is larger than, say, $1/2$ for all β large enough. \square

2 High-temperature expansion

As in the low-temperature case, the first step consists in proving an appropriate ‘geometrical’ representation of the sums involved in taking expectation with respect to the Ising model. To this purpose, given a subset $U \subseteq V$ of vertices, we let

$$Z_U(\beta) \equiv \sum_{x \in \{+1, -1\}^V} \prod_{i \in U} x_i \exp \left\{ \beta \sum_{(ij) \in E} x_i x_j \right\}. \quad (10)$$

Lemma 3. *For $U \subseteq V$, let $\mathfrak{G}(U)$ denote the set of subgraphs of G that have odd-degree in U and even degree (eventually vanishing) elsewhere. Then (denoting by $E(F)$ the set of edges of the subgraph F and by $|E(F)|$ its size)*

$$Z_U(\beta) = 2^{|V|} (\cosh \beta)^{|E|} \sum_{F \in \mathfrak{G}(U)} (\tanh \beta)^{|E(F)|}. \quad (11)$$

Proof Follows from the identity

$$\exp \left\{ \beta \sum_{(ij) \in E} x_i x_j \right\} = (\cosh \beta)^{|E|} \prod_{(i,j) \in E} (1 + x_i x_j \tanh \beta) \quad (12)$$

expanding the product on the right hand side and summing over x . \square

Lemma 4. *For $i, j \in V$, let $d(i, j)$ denote their graph theoretical distance. There exists $\beta_2 > 0$ such that, for all $\beta \leq \beta_2$, there exists constants $A(\beta), \lambda(\beta) \geq 0$ with $\lambda(\beta) \leq 1$, such that*

$$\mathbb{E}\{X_i X_j\} \leq A(\beta) \lambda(\beta)^{d(i,j)}. \quad (13)$$

Proof By the previous Lemma, and writing $t = \tanh \beta < 1$, we have

$$\mathbb{E}\{X_i X_j\} = \frac{\sum_{F \in \mathfrak{G}(\{i,j\})} t^{|E(F)|}}{\sum_{F' \in \mathfrak{G}(\emptyset)} t^{|E(F')|}}. \quad (14)$$

Notice that, for each connected component of the subgraph F , the sum of the degrees of its vertices is even. As a consequence, any $F \in \mathfrak{G}(\{i, j\})$ includes a path, call it F_{ij} , connecting i to j . Call the set of such paths $\mathfrak{P}(i, j)$. Further denote by $\mathfrak{G}(\emptyset; F_{ij})$ the set of subgraphs in $\mathfrak{G}(\emptyset)$ that have no edge in common with F_{ij} . Then we have

$$\mathbb{E}\{X_i X_j\} \leq \sum_{F_{ij} \in \mathfrak{P}(i,j)} t^{|E(F_{ij})|} \frac{\sum_{F' \in \mathfrak{G}(\emptyset; F_{ij})} t^{|E(F')|}}{\sum_{F' \in \mathfrak{G}(\emptyset)} t^{|E(F')|}} \leq \sum_{F_{ij} \in \mathfrak{P}(i,j)} t^{|E(F_{ij})|}. \quad (15)$$

The number of paths in $\mathfrak{P}(i, j)$ of length l is upper bounded by 3^l , and their minimal length $d(i, j)$. Therefore

$$\mathbb{E}\{X_i X_j\} \leq \sum_{l \geq d(i,j)} (3t)^l \leq (1 - 3t)^{-1} (3t)^{d(i,j)} \quad (16)$$

which proves the thesis. \square

Proof [Theorem 1, high temperature.] Notice that

$$\mathbb{E}\{\overline{X}^2\} = \frac{1}{N^2} \sum_{i,j \in V} \mathbb{E}\{X_i, X_j\} \leq \frac{1}{N^2} \sum_{i,j \in V} A \lambda^{d(i,j)} \leq \frac{A}{N} \sum_{j \in \mathbb{Z}^2} \lambda^{d(i,j)} \leq \frac{A}{N} \sum_{d=0}^{\infty} 8d \lambda^d \quad (17)$$

Since the last sum converges for $\lambda < 1$, we have $\mathbb{E}\{\overline{X}^2\} \leq C(\beta)/N$ for some constant $C(\beta)$ and all $\beta < \beta_2$, whence the thesis follows. \square

References

- [Fis67] Michael E. Fisher. Critical Temperatures of Anisotropic Ising Lattices. II. General Upper Bounds. *Phys. Rev.*, 162:480–485, 1967.
- [Gri64] Robert B. Griffiths. Peierls Proof of Spontaneous Magnetization in a Two-Dimensional Ising Ferromagnet. *Phys. Rev.*, 136:A437–A439, 1964.